What can thermal fluctuations tell us about fragility?

Thomas B. Schrøder

<u>Collaborators:</u> Ulf R.Pedersen Nicoletta Gnan Jon Papini Nicholas Bailey Jeppe C. Dyre

Fragility of Viscous Liquids: Cause(s) and Consequences, Copenhagen 2008



Danish National Research Foundation Centre for Viscous Liquid Dynamics

The single component Lennard-Jones liquid revisited - probably the most studied liquid in the history of computer simulations

Pressure and energy split in kinetic and <u>configurational</u> parts:

$$E(t) = K(t) + U(t) \qquad p(t)V = Nk_BT(t) + W(t)$$





[Pedersen et al. PRL 100 015701 2008] ²

Let's look at more state-points:

Each 'blob': scatter-plot of (W,U) over 10ns, after 10ns equilibration



3



The more subtle explanation

 $U_{LJ}(r) = kr^{-n} + br + U_0 + U_{rest}(r)$

- Explanation confirmed directly by simulations
- Effective exponent weakly dependent on state point
- Experimental data for supercritical Argon: R=0.96
- [N. Bailey et al. ArXiv:0807.055 (2008); to appear]

How general are the correlations?



- there exists a class of "strongly correlating liquids" [Pedersen et al. PRL 100 015701 2008]

Properties of strongly correlating viscous liquids, I

- Defining property: $\Delta W(t) = \gamma \Delta U(t)$, 'slope' slightly dependent on state point
- $\Delta p(t) = \frac{\gamma}{V} \Delta E(t)$, on "long time-scales"

$$p(t)V = Nk_BT(t) + W(t)$$
$$E(t) = K(t) + U(t)$$
$$+$$
Seperation of time-scales





[Pedersen et al. PRE 77 011201 2008]

Properties of strongly correlating viscous liquids, II

- Defining property: $\Delta W(t) = \gamma \Delta U(t)$, 'slope' slightly dependent on state point
- $\Delta p(t) = \frac{\gamma}{V} \Delta E(t)$, on "long time-scales"
- Fluctuation Dissipation theorem: $-Tc_{V}''(\omega) = \gamma^{2}K_{T}''(\omega) = -\gamma T \beta_{V}''(\omega)$ Three 'independent' thermoviscoelastic response functions are proportional. [Ellegaard et al., JCP 126, 074502 (2007); Pedersen et al., PRE 77, 011201 (2008)]
- Single-parameter aging: $W(t) = \gamma U(t) + W_0$ even out of equilibrium (isochoric!)





But: Is it the right form of scaling? What is the explanation?

Does not work for all viscous liquids



Density scaling: $\tau = F(\rho^{\gamma}/T)$

Suggested explanation: Effective power-law

See eg. [Coslovich & Roland, JPC 2008]

Conjectures:

- Density scaling if (and only if) strongly correlating liquid.
- "Density scaling exponent" = "fluctuation exponent"



Density scaling in the asymmetric dumbbell



Density scaling in LW-OTP



[Schrøder et al. ArXiv:0803.2199 (2008)]

Properties of strongly correlating viscous liquids, III

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- Single-parameter aging: $W(t) = \gamma U(t) + W_0$ even out of equilibrium (isochoric!)
- Density scaling (approximate): $au\,=\,F\,(
 ho^\gamma/T)$
- Scaling parameter from fluctuations (or pair of response-functions)
- Fragility: $m_P = m_\rho (1 + \alpha_P T_g \gamma)$ [K. Niss yesterday]

Thank you for your attention!



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1/12









Properties of strongly correlating viscous liquids, I • Defining property: $\Delta W(t) = \gamma \Delta U(t)$, 'slope' slightly dependent on state point • $\Delta p(t) = \frac{\gamma}{V} \Delta E(t)$, on "long time-scales" Data averaged over 0.1 au_{lpha} $\tau_{\alpha} \approx 1 \, ns$ 2 Normalized energy/pressure fluctuations (a) $p(t)V = Nk_BT(t) + W(t)$ Kob-Andersen BLJ E(t) = K(t) + U(t)0 -1 → Energy × ⊖ ← Pressure Seperation of time-scales Asymmetric dumbbells (b) 1 0 -1 Energy G - Pressure $\frac{3}{3}$ Time $[\tau_{\alpha}]$ 0 [Pedersen et al. PRE 77 011201 2008] 6

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/:

12